Luminescence in Lanthanum Chlorotungstate (LaWO₄Cl)

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The luminescence properties of LaWO₄Cl are reported and discussed. The tungstate group occurs as a trigonal bipyramid. Although there is some analogy with the luminescence of the tetrahedral tungstate group, the small Stokes shift reveals the presence of five oxygen ions. The luminescence of the following ions was studied: Mo^{6+} , Eu^{3+} , Sm^{3+} , and Tb^{3+} . The molybdate group shows red emission with a larger Stokes shift than the tungstate group. Energy transfer from the tungstate group to the rare earth ions occurs over nearest-neighbor distances, except for Tb^{3+} . Because of a low-lying f-d transition, the transfer efficiency is much higher for this ion.

1. Introduction

The luminescence of La₃WO₆Cl₃, and of several activators therein, has been reported recently (1, 2). The tungstate group in this lattice is trigonal prismatic (3). The crystal structure of LaWO₄Cl has also been described (4). The unusual feature of this structure is the five coordination of tungsten in the form of a trigonal bipyramid. It is the purpose of this paper to report on the luminescence of LaWO₄Cl and activators therein. In view of the similarity with the La₃WO₆Cl₃ system, the stress will be on the differences in the luminescence of both host lattices.

2. Experimental

The reader is referred to Refs. (1-4).

3. Results and Discussion

3.1. Unactivated LaWO₄Cl

The compound LaWO₄Cl does not luminesce at room temperature. At lower temperatures, however, a bright, blue luminescence appears under ultraviolet excitation. The quenching temperature is about 230K. Figure 1 presents the emission and excitation spectra at liquid helium temperature. The emission maximum is at 450 nm, the excitation maximum at 300 nm. This yields a Stokes shift of 11,000 cm^{-1} . This value is about equal to that reported for octahedral and trigonal prismatic tungstate groups (2). From this it can be concluded that Stokes shifts of 15,000 cm^{-1} and more in the case of tungstate emission are only to be expected for tetrahedral tungstate groups (2, 5).

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FIG. 1. Emission and excitation spectra of the luminescence of $LaWO_4Cl$ at LHeT.

Figure 2 shows the decay times of the WO_5^{4-} luminescence as a function of temperature. Above 4K the decay curves are exponential. Below 4K, they can be analyzed to give two decay times. The shorter one becomes too short at 4K to be measurable. This behavior differs from that observed for the octahedral and trigonal prismatic tungstate group, for which two radiant decay times were observed, each from a different energy level (2). The fast decay time at very low temperatures with its pronounced temperature dependence is ascribed to a nonradiative transition which feeds the emitting level. Since excitation is into the singlet levels and emission occurs from the triplet levels (5), this nonradiative transition may be the singlet-triplet transition.

The temperature dependence of the longer, radiative decay time is similar to

that reported for the tetrahedral tungstate group in CaWO₄ (6). It can be analyzed using a three-level scheme (5). The energy difference between the two excited states is found to be about 20 cm^{-1} . This value is also observed for CaWO₄. Unfortunately, no decay time measurements are known for CaWO₄ down to very low temperatures. Our results indicate that the emission in the case of LaWO₄Cl originates from a level which is connected with the ground state by a spin-forbidden, symmetry-allowed transition. The excited level is split by spin-orbit coupling. This splitting yields an optical trap level with a trap depth of $20 \text{ cm}^{-1}(5)$. It would be interesting to investigate whether the nonradiative, fast decay can also be observed in the case of CaWO₄. It has not been found for the WO_6^{6-} groups.

Vibrational spectra were measured to compare the W-O stretching frequencies of the WO₅⁴⁻ bipyramid with those of the tungstate tetrahedron and octahedron. The Raman spectrum was given in Ref. (4). In Table I, the infrared and Raman frequencies for LaWO₄Cl in the W-O frequency range are given. For an isolated WO₅⁴⁻ group, the symmetry should be D_{3h} . In LaWO₄Cl, the W⁶⁺ ion moves off the basal plane. Although the exact site symmetry is C_s , $C_{3\nu}$ is a good approximation. The W-O stretching frequencies have the following representations in this approximation: $3A_1 + E$. All of them are allowed in infrared and Raman.



FIG. 2. Decay times (τ) of the luminescence of LaWO₄Cl as a function of temperature. Also see text.

TABLE I

INFRARED AND RAMAN FREQUENCIES OF LaWO4Cl in the W-O STRETCHING REGION (ALL VALUES IN CM⁻¹)

Infrared ^a	Raman
950 (m) + 920 (sh)	925 (s)
850 (m)	825 (m)
760 (s)	765 (w)
730 (s)	730 (m)
710 (s)	715 (m)

^a Abbreviations: s = strong, m = medium, w = weak, sh = shoulder.

However, more than four lines are observed. Under C_s symmetry, the representations become 4A' + A''. The subcell described in Ref. (4) contains four W⁶⁺ ions, and has space group *Pnma* (D_{2h}^{16}). A detailed analysis becomes difficult. We assume that the five main frequencies (Table I) correspond to the five W-O frequencies expected for C_s site symmetry. The two highest frequencies do not coincide in infrared and Raman, which indicates factor group splitting.

The values of the W–O frequencies of LaWO₄Cl are in the same region as those for CaWO₄ (7). Those for the WO₆⁻ groups are considerably lower (1, 8).

Although the crystallographic and vibrational spectroscopy data suggest a close analogy between the W-coordination in La WO_4Cl and that in a tungstate tetrahedron, it seems to us that the value of so sensitive a parameter as the Stokes shift clearly indicates the presence of a fifth oxygen ion. Probably the W^{6+} ion is more off-center in the ground state than in the excited state.

3.2. $LaW_{1-x}Mo_xO_4Cl$

Samples with a few percent of molybdenum could be prepared easily. Emission and excitation spectra are presented in Fig. 3. The molybdate emission peaks at about 650 nm. The excitation spectrum shows two bands, viz., one at about 330 nm corresponding to molybdate absorption and one at 300 nm corresponding to tungstate absorption. The latter indicates energy transfer from tungstate to molybdate, but because of its low intensity this transfer is of limited importance.

The Stokes shift of the molybdate emission of LaWO₄Cl:Mo is about 14,500 cm⁻¹, which is much larger than that for the tungstate emission (see above). Usually the Stokes shift of similar molybdate and tungstate groups has the same value, e.g., for CaWO₄ and CaMoO₄ (9). The larger Stokes shift for LaWO₄Cl:Mo may well indicate that the smaller Mo⁶⁺ ion has moved even further away from the basal plane of the pyramid than the W⁶⁺ ion, so that a tetrahedral coordination occurs.

The quenching temperature of the molybdate luminescence is 140K. In contradistinction with the molybdate emission in La₃ WO_6Cl_3 , that in LaWO₄Cl has a high



FIG. 3. Emission and excitation spectra of the molybdate luminescence of LaW_{0.97}Mo_{0.03}O_4Cl at LHeT.

quantum efficiency at low temperatures. This made it possible to perform decay time measurements. Figure 4 presents the decay time as a function of temperature. There is a maximum at 20K. At higher temperatures the decay time decreases as a result of thermal quenching. At lower temperatures, the decay time goes to a value of 160 µsec. Note that the values at 20K and higher are 1 order of magnitude larger than for the corresponding tungstate luminescence. This phenomenon was also observed for the tetrahedral groups and is due to a weaker spin-orbit coupling in the case of molybdenum. The weaker coupling makes the triplet-singlet emission transition more forbidden.

The level splitting responsible for the low-temperature behavior is different in the molybdate and the tungstate. This splitting is due to spin-orbit effects in the latter case (5). The energy difference between the two levels involved is estimated to be 5 cm⁻¹, which agrees reasonably well with the value for CaMoO₄ (5). For details of analysis and interpretation, the reader is referred to a recent review paper (5).

It is hard to explain why the red emission of La₃WO₆Cl₃:Mo (2) has a low quantum efficiency at LHeT, whereas that of La WO₄Cl:Mo does not. For an explanation, it is necessary to know the complicated level structure of the excited state. Note, in addition, that the quenching temperature of the luminescence of La₃WO₆Cl₃ is higher than



FIG. 4. Decay times of the molybdate luminescence of $LaW_{0.97}Mo_{0.03}O_4Cl$ as a function of temperature.

that of LaWO₄Cl, which makes the problem even more difficult.

3.3. LaWO₄Cl-Sm

Several rare earth ions show luminescence in LaWO₄Cl (4). Here the luminescence properties are reported in more detail, stressing the energy transfer from host lattice to the rare earth ions.

First, we consider the Sm³⁺ ion. Results presented here are for the composition La_{0.98}Sm_{0.02}WO₄Cl. At room temperature, the characteristic Sm³⁺ emission can be excited in the Sm^{3+} lines (4). The excitation spectrum at 300K does not show the tungstate absorption band at 300 nm, so that at 300K energy transfer from the WO_5^{4-} group to the Sm³⁺ ion does not occur. At LHeT. the situation is different. Upon tungstate excitation a complicated emission spectrum is observed consisting of the characteristic Sm³⁺ lines (for about 12%) and the tungstate emission band (for about 88%). The tungstate emission band shows a number of sharp dips coinciding with the sharp Sm³⁺ excitation lines. From this, we conclude that nonradiative as well as radiative energy transfer occurs. An analysis is given in Ref. (2) and shows that the nonradiative transfer occurs only over nearest-neighbor distances. Each tungstate group has six nearest-neighbor La³⁺ sites. For a Sm³⁺ concentration of 2 at%, this means that $0.98^6 = 89\%$ of the tungstate groups have no Sm³⁺ neighbors and decay radiatively. However, 11% of the tungstate groups see one or more Sm³⁺ ions and transfer their excitation energy, so that Sm³⁺ emission occurs. From the emission spectrum, it is estimated that about 1% of the tungstate emission is, in addition, transferred radiatively. In this model, the Sm³⁺-tungstate emission ratio is predicted to be 0.14, in excellent agreement with the experimental value.

It is interesting to follow the emission spectrum under tungstate excitation as a

function of temperature. Up to 100K nothing changes, but at higher temperatures the tungstate emission band loses intensity. Above 230K, it has disappeared. Also the Sm^{3+} emission intensity decreases, but not as strongly as the tungstate emission intensity. At 300K, about 30% of the low-temperature intensity is left.

From this, we can conclude the following: (1) energy migration among the tungstate groups in LaWO₄Cl does not occur. A similar observation was made for La₃ WO₆Cl₃ (2). At higher temperatures the nonradiative processes in the tungstate group are so fast that they compete with the WO₅⁴⁻ \rightarrow Sm³⁺ transfer probability.

3.4. $LaWO_4Cl:Eu^{3+}$

Samples with Eu³⁺ concentrations below 1 at% did not show Eu³⁺ luminescence. A sample with 5% Eu³⁺, i.e., La_{0.95} Eu_{0.05}WO₄Cl, showed efficient Eu³⁺ emission and was studied in more detail. At room temperature, the results are similar to those for Sm³⁺, i.e., only Eu³⁺ emission and excitation lines were observed in the spectra. At LHeT, the emission spectrum under Eu³⁺ excitation shows a complete splitting of all the emission transitions in accordance with the low site symmetry (C_s), i.e., ${}^5D_0{}-{}^7F_0$ 1, ${}^5D_0{}-{}^7F_1$ 3, ${}^5D_0{}-{}^7F_2$ 5, etc. (Fig. 5).

Upon tungstate excitation at LHeT, this sample shows mainly tungstate emission, viz., 94% of the total emission. Only 6% of the emission originates from the Eu³⁺ ions. A small amount of radiative transfer could just be observed. In the simple model described above for LaWO₄Cl:Sm³⁺, we predict 26% of Eu³⁺ emission. Since it does not seem realistic to assume an inefficient energy transfer from tungstate to Eu³⁺, we conclude that the real concentration of Eu³⁺ ions is less than 5 at%. This was also observed for La₃WO₆Cl₃:Eu³⁺ (2) and has been ascribed to a slight amount of reduction in the preparation procedure. It may be



FIG. 5. Emission spectrum of $La_{0.95}Eu_{0.05}WO_4Cl$ under Eu^{3+} excitation at LHeT in the ${}^5D_{0-}{}^2F_{0,1,2}$ region.

that the absence of Eu^{3+} emission in the lower-concentration samples is due to the same reason. An amount of Eu^{3+} emission of 6% points to a concentration of 1 at% using our energy-transfer description.

As in the case of La₃WO₆Cl₃:Eu³⁺, the charge-transfer band of Eu³⁺ in LaWO₄ Cl:Eu³⁺ is at relatively low energy. Figure 6 shows the excitation spectra of the Eu³⁺ and the tungstate emission of LaWO₄ Cl:Eu³⁺. The extra broadband on the low-energy side of the tungstate band is the Eu³⁺ charge-transfer band. This phenomenon is similar to that observed for La₃WO₆Cl₃:Eu³⁺, where the charge-transfer band is at even lower energy.

3.5. LaWO₄Cl:Tb³⁺

For Tb^{3+} activation, an efficient, green phosphor is obtained. It turns out that there are striking differences to Sm^{3+} and Eu^{3+} . Data given here refer to the composition $La_{0.98}Tb_{0.02}WO_4Cl$. At room temperature, there is efficient, green emission. The excitation spectrum consists not only of the characteristic Tb^{3+} lines, but also of a broadband which coincides partly with the tungstate excitation band. This shows that at room temperature energy transfer from



FIG. 6. The excitation spectra of the Eu^{3+} and the tungstate emission of $La_{0.95}Eu_{0.05}WO_4Cl$ at LHeT. Note the additional broadband in the case of Eu^{3+} , which corresponds to the charge-transfer transition.

tungstate to Tb^{3+} occurs, whereas it does not to Sm^{3+} and Eu^{3+} .

At LHeT, the emission spectrum under tungstate excitation consists of a weak tungstate emission band (50% of the total emission intensity) and sharp Tb^{3+} lines (${}^{5}D_{4-}{}^{7}F_{J}$), which also contribute 50%. The total emission intensity is as strong as in the case of Sm³⁺- and Eu³⁺-activated samples. For the Tb³⁺ concentration concerned, we expect the same Tb³⁺-tungstate emission ratio as derived above for LaWO₄Cl:Sm³⁺, i.e., 0.14. The experimental value, viz., 1.0, is much higher. From this, we have to conclude that the tungstate-Tb³⁺ transfer proceeds over considerably longer distances than those between nearest neighbors.

The reason for this may become clear from Fig. 7, which gives the excitation spectrum of the Tb^{3+} emission at LHeT. Note the broad tail on the low-energy side

of the tungstate excitation band. This tail is due to direct Tb^{3+} excitation. We have to ascribe it to one of the components of the $4f \rightarrow 5d$ transitions on the Tb^{3+} ion. This is an allowed optical transition. A comparison between Figs. 1 and 7 shows that there is spectral overlap between the tungstate emission and the broad Tb^{3+} tail, i.e., the tungstate emission overlaps an allowed Tb^{3+} absorption band. In the case of Sm^{3+} and Eu^{3+} , there is only overlap with forbidden 4f-4f transitions. This gives the tungstate- Tb^{3+} transfer a considerably higher range. A similar situation exists in YTa O_4 :Tb (10).

In the LaWO₄Cl structure, the La³⁺ ions are situated on rows along the *b* axis (Ref. (4), Fig. 2). A 5*d* orbital oriented along this direction will have a very low energy because of the attractive forces of the La³⁺ ions. This seems to be an obvious explana-



FIG. 7. The excitation spectrum of the Tb^{3+} emission of La_{0.98}Tb_{0.02}WO₄Cl at LHeT. Sharp peaks are Tb^{3+} *f*-*f* transitions, the band peaking at 300 nm is the tungstate excitation band, and the broad tail from 390 up to 300 nm is the Tb^{3+} *f*-*d* transition.

tion for the exceptionally low position of one of the 4f-5d components in the spectra.

3.6. Other Ions

The U⁶⁺ as well as the Bi³⁺ ion were found to luminesce efficiently in La₃WO₆Cl₃ (1, 2). Therefore, we also tried these ions in LaWO₄Cl, but no emission of any intensity could be observed.

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